Comparison of Numerical Solution of 1D Hyperbolic Telegraph Equation using B-Spline and Trigonometric B-Spline by Differential Quadrature Method

Geeta Arora and Varun Joshi*

Department of Mathematics, Lovely Professional University, Punjab – 144411, India; geetadma@gmail.com, varunjoshi20@yahoo.com

Abstract

Objectives: This paper aims to compute the approximate solution of one dimensional (1D) hyperbolic telegraph equation with appropriate primary and limiting conditions. **Methods/Statistical Analysis**: To find the approximate solution, two different modified spline basis function are used with the differential quadrature method and splines are used to compute the weighting coefficients and thus the equation is transformed to a set of first order conventional differential equations which is further solved by the SSP-RK43 method. Three test problems of this equation are simulated to establish the precision and usefulness of the proposed scheme. **Findings:** The obtained numerical results are found to be good in terms of accuracy, efficiency and simplicity. To validate the computed results using proposed scheme, various comparisons at different time levels has been done in the form of L_2 and L_{∞} errors. These errors are compared and enlisted in the form of tables with computed errors enlisted in literature. **Application/Improvements:** Being an important equation of nuclear material science, one dimensional (1D) hyperbolic telegraphs equation needs to take care in the sense of better numerical solution. In this context, a successful effort has been done in this research article by proposing a hybrid numerical scheme.

Keywords: 35K57, Differential Quadrature Method, Hyperbolic Telegraphs Equation, Mathematics Subject Classification (2010): 65M06, Modified Cubic B-Spline, Modified Trigonometric B-Spline, SSP-RK43

1. Introduction

The telegraph equation is utilized to outline the reaction diffusion in numerous branches of emerging sciences. This mathematical equation premised for crucial equations of nuclear material science. It is used to demonstrate the vibrations of structures, e.g. structures, shafts, and machines. It likewise emerges in the investigation of throb blood stream in supply routes, in the 1D irregular movement of bugs along a hedge^{1,2} and also play a significant part in demonstrating of numerous appropriate problems like signal investigation³, wave propagation⁴, random walk theory⁵ etc. This mathematical equation is regularly utilized as a part of signal investigation for transmission

 $u_t(x,t) + 2\alpha u_t(x,t) + \beta^2 u(x,t) = u_x(x,t) + f(x,t) \quad x \in [a,b] \quad t \ge 0$ (1)
With initial or starting conditions (ICs)

$$u(x,t_0) = g_1(x) , u_t(x,t_0) = g_2(x), x \in [a,b],$$

and with limiting conditions (BCs) (2)

$$u(a,t) = \psi_0$$
, $u(b,t) = \psi_1$, $t > 0$ (3)

where $f, g_1, g_2, \psi_1, \psi_2$ are known functions

and
$$u_{tt} = \frac{\partial^2 u}{\partial t^2}, u_t = \frac{\partial u}{\partial t}, u_{xx} = \frac{\partial^2 u}{\partial x^2}.$$

and proliferation of electrical signs⁶ furthermore has applications in different fields⁷. This equation is given as:

^{*}Author for correspondence

Eq. (1) with the coefficients $\beta = 0$, $\alpha > 0$ stand for a partial differential equation which is actually a damped wave equation and correspond to the telegraph mathematical equation if $\alpha > 0$, $\beta > 0$. The telegraph equation is likewise adverted to show the combination of dissemination and wave proliferation by the properties of constrained velocity to standard warmth or mass transport condition¹.

In past few years, many numerical schemes to solve telegraph equation have been developed by many researchers. The existence of its double periodic solution was investigated in⁸. The unconditionally stable schemes for telegraph Eq. (1) was used in⁹⁻¹¹, etc. Further, different conditionally stable finite difference schemes were implemented to solve the telegraph equation¹²⁻¹⁴. In¹⁵ finite difference for approximation of spatial derivatives and the time derivative handled by collocation scheme is used. Chebyshev cardinal functions were used for numerical simulation of the equation in¹⁶. Various other numerical schemes were also developed or implemented to solve the telegraph equation, Some of them are: method using interpolating scaling functions¹⁷, Chebyshev Tau method¹⁸, Rothe-wavelet method¹, semi-discretion methods¹⁹, explicit difference method¹⁴, DRBIE method²⁰, collocation scheme along thin plate splines radial basis function (RBFCM)²¹, Polynomial Differential Quadrature Method (PDQM)²², Cubic B-spline Quasi-interpolation (CBQ)²³, cubic B-spline collocation method (CBCM)²⁴, Quartic B-spline Collocation Method (QBCM)²⁵, and Collocation Method based on Modified Cubic B-spline (CMMCB)²⁶, etc.

This article is anxious with an approximate solution of 1D hyperbolic PDE with suitable starting and limiting conditions acquired by utilizing altered cubic B-spline and trigonometric B-spline with differential quadrature strategy, consequently alluded as (MCB-DQM) and (MTCB-DQM)²⁷. The proposed method depends on differential quadrature technique with cubic-B-spline and trigonometric B-spline basis function. Eq. (1) is initially changed over into an arrangement of coupled Partial Differential Equations (PDEs) utilizing the transformation: $u_t = v$. The MCB-DQM and MTCB-DQM are applied to change the PDEs into a system of first order ODEs, in time, which is solved by utilizing the SSP-RK43 scheme²⁸. Three test problems of this equation have been considered to exhibit the exactness and utility of the proposed technique. The maximum absolute error L_{∞} and L_{2} error in the MCB-DQM and MTCB-DQM

approximate solution have been analogized with the error calculated by other numerical schemes accessible in the literature. The numerical results are seen to be in awesome simultaneousness with the exact solution.

The article is partitioned into various sections, in Section 2, the description of the modified cubic B-spline and modified trigonometric B-spline with Differential Quadrature Method (DQM) is given. Sections 3, discusses the procedure to utilize the proposed scheme. Three numerical test problems are given in to fabricate the significance and precision of the proposed technique in Section 4. In last section 5, summery of the paper is given in form of conclusion.

2. Description of the Method

The DQM was initially proposed by²⁹. This strategy has been effectively connected to understand different one and two dimensional differential equations by utilizing various basis functions³⁰⁻³³.

In DQM the derivatives of some function which are present in the given PDE are replaced with their approximate values at desired different points. Because of the dependence on weight coefficients to domain grid points, we consider N equidistance nodal points on the real axis, that is $a = x_1 < x_2 < \cdots < x_{N-1} < x_N = b$ with $x_{i+1} - x_i = h$. The solution u(x, t) at knot x_i Is denoted by $u(x_i, t)$ for $i = 1, 2, 3, \dots, N$. The approximate values of spatial derivatives are given as:

$$u_{x}(x_{i},t) = \sum_{j=1}^{N} a_{ij} u(x_{j},t), \qquad u_{xx}(x_{i},t) = \sum_{j=1}^{N} b_{ij} u(x_{j},t), \qquad i = 1, 2, \dots, N$$
(4)

2.1 Differential Quadrature Method uses B-Spline

As B-spline have the certain nice properties like smoothness and competence to handle indigenous singularities, as B-spline basis functions are easy to implement so many researchers used cubic B-spline basis function to compute the approximate solution of physical models^{26,27} definedas follows:

$$\varphi_{\downarrow}j(x) = 1/h^{\dagger}3 \left\{ \bullet ((x - x_{\downarrow}(j - 2))^{\dagger}3, x \in [x_{\downarrow}(j - 2, ...) \cap x \cap_{\downarrow}(j - 1)) \otimes (x - x_{\downarrow}(j - 2))^{\dagger}3 - 4(x - x_{\downarrow}(j - 1))^{\dagger}3, x \in [x_{\downarrow}(j - 1, ...) \cap x \cap_{\downarrow}j) \otimes \bullet ((x_{\downarrow}(j + 2) - x)^{\dagger}3 - 4(x_{\downarrow}(j - 1))^{\dagger}3, - 4(x_{\downarrow}(j - 1))^{\dagger}3, x \in [x_{\downarrow}(j - 1, ...) \cap x \cap_{\downarrow}j) \otimes \bullet ((x_{\downarrow}(j + 2) - x)^{\dagger}3 - 4(x_{\downarrow}(j - 1))^{\dagger}3, - 4(x_{\downarrow}(j - 1))^{\dagger}3, x \in [x_{\downarrow}(j - 1, ...) \cap x \cap_{\downarrow}j) \otimes \bullet ((x_{\downarrow}(j - 2) - x)^{\dagger}3 - 4(x_{\downarrow}(j - 1))^{\dagger}3, - 4(x_{\downarrow}(j - 1))^{\dagger}3, x \in [x_{\downarrow}(j - 1, ...) \cap x \cap_{\downarrow}j) \otimes \bullet ((x_{\downarrow}(j - 2) - x)^{\dagger}3 - 4(x_{\downarrow}(j - 1))^{\dagger}3, x \in [x_{\downarrow}(j - 1, ...) \cap x \cap_{\downarrow}j) \otimes \bullet ((x_{\downarrow}(j - 2) - x)^{\dagger}3 - 4(x_{\downarrow}(j - 1))^{\dagger}3, x \in [x_{\downarrow}(j - 1) \cap x \cap_{\downarrow}j) \otimes \bullet ((x_{\downarrow}(j - 2) - x)^{\dagger}3 - 4(x_{\downarrow}(j - 1))^{\dagger}3, x \in [x_{\downarrow}(j - 1) \cap x \cap_{\downarrow}j) \otimes \bullet ((x_{\downarrow}(j - 2) - x)^{\dagger}3 - 4(x_{\downarrow}(j - 1))^{\dagger}3, x \in [x_{\downarrow}(j - 1) \cap x \cap_{\downarrow}j) \otimes \bullet ((x_{\downarrow}(j - 2) - x)^{\dagger}3 - 4(x_{\downarrow}(j - 2))^{\dagger}3 - 4(x_{\downarrow}($$

where $\{\varphi_0, \varphi_1, \dots, \varphi_N, \varphi_{N+1}\}$ generate the basis on [a, b].

Lemma 1: The numerical values of φ_i and its derivatives φ'_i, φ''_i at j th nodal point are evaluated as

$$\varphi_{i}(x_{j}) = \begin{cases} 4, & \text{if } i-j = \mathbf{0} \\ 1, & \text{if } i-j = \pm 1, \varphi_{i}'(x_{j}) = \begin{cases} \pm \frac{3}{h}, & i-j = \pm 1, \\ 0, & \text{else} \end{cases},\\ \varphi_{i}''(x_{j}) = \begin{cases} -\frac{12}{h^{2}}, & \text{if } i-j = \mathbf{0} \\ \frac{6}{h^{2}}, & \text{if } i-j = \pm 1 \\ 0, & \text{else} \end{cases}$$
and

The first order derivative approximation at the grid point x_i , i = 1, 2, ..., N is given by

$$\phi'_{k}(x_{i}) = \sum_{j=1}^{N} a_{ij} \phi_{k}(x_{j}), \qquad k = 1, 2, ..., N.$$
(6)

By Lemma 1 and modified basis function mentioned in section 2.2, Eq. (6) is reduced to linear equations as

$$A\vec{a}[i] = \vec{H}[i], \text{ for } i = 1, 2, ..., N.$$
 (7)

and A represented as:

$$A = \begin{bmatrix} 6 & 1 & & & \\ 0 & 4 & 1 & & \\ & 1 & 4 & 1 & & \\ & & \ddots & \ddots & \ddots & \\ & & 1 & 4 & 1 & \\ & & & 1 & 4 & 0 \\ & & & & 1 & 6 \end{bmatrix}$$

weighting coefficient vector with respect to the points x_i , are represented as $\vec{a}[i]$, that is

 $\vec{a}[i] = [a_{i1}, a_{i2}, a_{i3}, \dots a_{iN}]^T$ and the coefficient vector $\vec{h}[i] = [h_{i1}, h_{i2}, h_{i3}, \dots h_{iN}]^T$

vector $\vec{h}[i] = [h_{i1}, h_{i2}, h_{i3}, \dots h_{iN}]^T$ with respect to x_i , $i = 1, 2, \dots, N$ are evaluated as:

$$\vec{H}[\mathbf{1}] = \begin{bmatrix} -\frac{6}{h} \\ \frac{6}{h} \\ 0 \\ \vdots \\ 0 \\ 0 \end{bmatrix}, \vec{H}[\mathbf{2}] = \begin{bmatrix} -\frac{3}{h} \\ 0 \\ \frac{3}{h} \\ \vdots \\ 0 \\ 0 \end{bmatrix}, \dots, \vec{H}[N-\mathbf{1}] = \begin{bmatrix} 0 \\ \vdots \\ 0 \\ -\frac{3}{h} \\ 0 \\ \frac{3}{h} \end{bmatrix}, \qquad \vec{H}[N] = \begin{bmatrix} 0 \\ 0 \\ \vdots \\ 0 \\ -\frac{6}{h} \\ \frac{6}{h} \end{bmatrix}$$

2.2 Differential Quadrature Method using Trigonometric B-Spline

The cubic trigonometric B-spline basis function $TB_m(x)$ for m = -1, 0, 1, ..., N + 1 is characterized as:

where
$$p(x_m) = \frac{\sin((\mathbf{x} - \mathbf{x}_m))}{2}$$

$$q(x_m) = \frac{\sin((\mathbf{x_m} - \mathbf{x}))}{2}, \text{ and}$$
$$w = \sin\left(\frac{h}{2}\right)\sin(h)\sin\left(\frac{3h}{2}\right) \text{ and } h = \frac{b-a}{n} \tag{8}$$

and $TB_m(x)$ is cubic trigonometric B-spline basis function with some geometric properties like

 C^{∞} Continuity, non-negativity and partition of unity. The values of $TB_m(x)$, $TB'_m(x)$ and $TB''_m(x)$ are given by Table 1, where:

$$a_{1} = \frac{\sin^{2}{h}}{\sin(h)\sin(\frac{2h}{2})}, a_{2} = \frac{2}{1+2\cos(h)}, a_{3} = -\frac{3}{4\sin(\frac{2h}{3})}, a_{4} = \frac{3}{4\sin(\frac{2h}{3})}, a_{5} = \frac{3(1+3\cos(h))}{16\sin^{2}{h}(2\cos(\frac{2h}{3}) + \cos(\frac{2h}{3}))}, and$$

$$a_{6} = -\frac{3\cos^{2}{h}(\frac{h}{2})}{\sin^{2}{h}(\frac{h}{2})(2+4\cos(h))}$$

Both above defined basis functions can be improved so that the resulting matrix becomes diagonally dominant. The modification of these functions can be done as follows: 2Z

$$B_{1}(x) = B_{1}(x) + 2B_{0}(x),$$

$$B_{2}(x) = B_{2}(x) - B_{0}(x),$$

$$B_{j}(x) = B_{j}(x) \text{ for } j = 3, 4, ..., N - 2$$

$$B_{N}(x) = B_{N}(x) + 2B_{N+1}(x),$$

$$B_{N-1}(x) = B_{N-1}(x) - B_{N+1}(x),$$

where{ $B_{1}, B_{2}, ..., B_{N}$ } form a basis on [a, b]
For trigonometric B-spline the system becomes

For trigonometric B-spline the system becomes $A\vec{a}[i] = \vec{R}[i], \text{ for } i = 1, 2, ..., N.$

where A A represented as:

$$\begin{bmatrix} a_{2} + 2a_{1} & a_{1} & 0 & 0 & . & 0 \\ 0 & a_{2} & a_{1} & . & . & 0 \\ 0 & a_{1} & a_{2} & a_{1} & . & 0 \\ 0 & 0 & a_{1} & a_{2} & a_{1} & 0 \\ 0 & . & . & . & 0 \\ 0 & 0 & . & . & a_{1} & a_{2} + 2a_{1} \end{bmatrix}$$
(9)

Weighting coefficient vector with respect to the points x_i , are represented as $\vec{a}[i]$, that is

 $\vec{a}[i] = [a_{i1}, a_{i2}, a_{i3}, \dots, a_{iN}]^T$ and the coefficient vector $\vec{R}[i] = [r_{i1}, r_{i2}, r_{i3}, \dots, r_{iN}]^T$

Table 1. The values of TB_m , TB'_m and TB''_m at different node points

TB\x	x_{m-2}	x_{m-1}	x_m	x_{m+1}	x_{m+2}
TB_m	0	<i>a</i> ₁	<i>a</i> ₂	<i>a</i> ₁	0
TB_{m}^{\prime}	0	<i>a</i> 3	0	a_4	0
$TB_m^{\prime\prime}$	0	a 5	а ₆	a 5	0

with respect to x_i , i = 1, 2, ..., N, or can be presented as:

$$\vec{R}[\mathbf{1}] = [-2a_4, a_3 - a_4, \quad 0, \dots, \mathbf{0}]^T,$$

$$\vec{R}[\mathbf{2}] = [a_4, 0, a_3, \quad 0, \dots, \mathbf{0}]^T,$$

 $\vec{R}[\mathbf{3}] = [0, \qquad a_{\mathbf{4}}, 0, a_{\mathbf{2}}, \qquad 0, \dots, \mathbf{0}]^T, \dots, \vec{R}[N-1] = [0, \dots, \qquad 0, a_{\mathbf{4}}, 0, a_{\mathbf{3}}]^T, \qquad \vec{R}[N] = [0, \dots, 0, \qquad a_{\mathbf{4}} - a_{\mathbf{2}}, 2a_{\mathbf{2}}]^T$

Now we apply the well-known "Thomas Algorithm" to compute the solution of the obtained equations which provides the vector $\vec{a}[i]$, using the coefficient a_{ij} , the weighting coefficients b_{ij} , for i = 1, 2, 3, ..., N, j = 1, 2, 3, ..., N are evaluated as follows:²⁸ $b_{ij} = 2a_{ij} \left(a_{ii} - \frac{1}{x_i - x_j} \right)$, for $i \neq j$, and $b_{ii} = -\sum_{i=1,i\neq j}^{N} b_{ij}$

3. Implementation of the Method

Using the transformation: $u_t = v$ the Eq. (1) is transformed to a set of PDEs as follows:

$$u_t(x,t) = v(x,t)$$

$$v_t(x,t) = u_{xx}(x,t) - 2\alpha v(x,t) - \beta^2 u(x,t) + f(x,t) \quad (10)$$

On substituting the values for the second order approximation of the space derivatives, obtained from MCB-DQM and MTCB-DQM, Eq. (10) rewritten as: u(x, t) = v(x, t)

$$v_{t}(x_{i},t) = v(x_{i},t)$$

$$v_{t}(x_{i},t) = \sum_{j=1}^{N} b_{ij}u(x_{j},t) - 2\alpha v(x_{i},t) - \beta^{2}u(x_{i},t) + f(x_{i},t) \quad (11)$$
for $i = 1, 2, ..., N$

Hence, Eq. (11) reduces into a coupled system of firstorder ODEs in time, that is,

$$\frac{du_i}{dt} = v_i \text{ and}$$

$$\frac{dv_i}{dt} = [L(u]_i), \quad i = \{1, 2, \dots, N\}, \quad (12)$$

where L represents the right hand side of ODEs.

The equation is solved subject to the BCs as defined in (3), and the ICs.

$$u(x_i, \mathbf{0}) = g_i(x_i), \qquad v(x_i, \mathbf{0}) = g_2(x_i), \qquad i \in \{1, 2, \dots, N\},$$
(13)

The resulting set of ODEs are solved by using SSP-RK43 scheme²⁸ given below:

$$\begin{aligned} u^{(1)} &= u^m + \frac{\Delta t}{2} L(u^m) \\ u^{(2)} &= u^{(1)} + \frac{\Delta t}{2} L\left(u^{(1)}\right) \\ u^{(3)} &= \frac{2}{3} u^m + \frac{u^2}{3} + \frac{\Delta t}{6} L\left(u^{(2)}\right) \\ u^{m+1} &= u^{(3)} + \frac{\Delta t}{2} L\left(u^{(3)}\right) \end{aligned}$$

and consequently the solutions u(x, t), at the required time level are obtained.

4. Numerical Experiments

For numerical discussion, three test problems are considered to obtain the approximate solutions by MCB-DQM and MTCB-DQM. The L_2 and L_{∞} error norms are calculated using the exact solution and formulas to calculate these errors are defined in Eq. (14).

$$L_{2} = \left(h \sum_{j=1}^{N} \left[u_{j}^{exact} - u_{j}^{*}\right]^{2}\right)^{\frac{1}{2}} \text{ and}$$
$$L_{\infty} = \max_{\substack{j=1\\j=1}}^{N} \left|u_{j}^{exact} - u_{j}^{*}\right|, \qquad (14)$$

where u_j represent the numerical solution at node j.

Example 1: The telegraph equation over the region $[0, \pi]$ is considered with following conditions

$$u(0, t) = 0, \qquad u(\pi, t) = 0, \qquad t \ge 0$$

and the function

$$f(x,t) = (2 - 2\alpha + \beta^2)e^{-t}\sin(x)$$

The exact solution^{19,20} is given by:

$$u(x,t) = e^{-t} \sin(x) \tag{15}$$

The comparison of L_2 and L_{∞} errors at different time levels $t \leq 2$ is done for $\alpha = 2$, $\beta = \sqrt{2}$, $\alpha = 3$, $\beta = \sqrt{2}$ (a) For $\alpha = 2$, $\beta = \sqrt{2}$ the L_2 and L_{∞} errors with h = 0.02, $\Delta t = 0.01$ are compared with the errors due to the well-known earlier schemes: CMMCB²⁶ and RBFCM²¹, and are reported in Table 2. It is evident that the errors are decreasing with increment in time, and the numerical results are more accurate than numerical solutions obtained by CMMCB²⁶ and RBFCM²¹. The physical conduct of solutions obtained by MCB-DQM and MTCB-DQM at various time levels $t \leq 2$ is depicted in Figure 1.

Table 2. Comparison of L_2 and L_{∞} errors at $t \le 2$ with h = 0.02 and $\Delta t = 0.01$.

t	L_2			
	MTCB - DQM	MCB - DQM	CMMCB ²⁶	RBFCM ²¹
0.5	1.4022E - 07	2.1570E - 07	2.3328E - 06	7.9491E - 05
1.0	1.059E - 07	1.6491E - 07	4.3667E - 06	1.4554E - 04
1.5	7.2312E - 08	1.1434E - 07	4.7817 E - 06	1.5895E - 04
2.0	5.0662E - 08	8.0521E - 08	4.2706 E - 06	1.4185E - 04
	L_{∞}			
0.5	1.9891E - 07	3.5160E - 07	1.8612E - 06	8.3721E - 06
1.0	1.1694E - 07	2.1430E - 07	3.4839E - 06	1.5680E - 05
1.5	7.1148E - 08	1.3045E - 08	3.8251E - 06	1.7412E - 05
2.0	4.3259E - 08	7.9350E - 08	3.4073E - 06	1.5813E - 05



Figure 1. At different time levels $t \le 2$ the performances of exact (Left) and approximate solution of Example 1 with $\alpha = 2$, $\beta = \sqrt{2}$ using trigonometric B-spline (Middle) and B-spline (Right).

(b) For $\alpha = 3$, $\beta = \sqrt{2}$ the L_2 and L_{∞} errors with h = 0.02, $\Delta t = 0.0001$ are compared with the errors by RBFCM²¹, and are reported in Table 3. It is evident that the errors are decreasing as time increases (also, by decreasing the values of h), and computed results are more accurate than the results obtained in²¹. The physical behaviour of the MCB-DQM and MTCB-DQM solutions for $\Delta t = 0.0001$ are depicted in Figure 2 at various time levels $t \leq 2$.

Table 3. Compari	son of L 2	and L ∞	errors at t	≤ 2
with $h = 0.02$	and $\Delta t =$	0.000	1	

t	L ₂		
	MTCB - DQM	MCB - DQM	RBFCM ²¹
0.5	5.5700E - 06	6.22E - 05	1.55E - 05
1.0	9.7571E - 06	1.05E - 05	2.71E - 05
1.5	1.0903E - 05	1.16E - 05	3.03E - 05
2.0	1.0407E - 05	1.11E - 05	2.92E - 05
	L_{∞}		
0.5	4.4591E - 06	4.46E - 06	1.41E - 04
1.0	7.8127E - 06	7.81E - 06	2.46E - 04
1.5	8.7306E - 06	8.74E - 06	2.74E - 04
2.0	8.3216E - 06	8.53E - 06	2.62E - 04



Figure 2. At different time levels $t \le 2$ the performances of exact (Left) and approximate solution of Example 2 with $\alpha = 3$, $\beta = \sqrt{2}$ using trigonometric B-spline (Middle) and B-spline (Right).

Example 2: The telegraph equation is considered over the region [0, 1] for $\alpha = 0.5$, $\beta = 1$ with conditions.

 $u(x, \mathbf{0}) = 0,$ $u_t(x, \mathbf{0}) = 0,$ u(0, t) = 0, u(1, t) = 0, $t \ge \mathbf{0}$

and the function

 $f(x, t) = (2 - 2t + t^2)(x - x^2)e^{-t} + 2t^2e^{-t}$ The analytic solution of the equation^{18,21} is given as: $u(x, t) = (x - x^2)t^2e^{-t}$



Figure 3. At different time levels $t \le 5$ the performances of exact (Left) and approximate solution of Example 2 with $\Delta t = 0.0001$, h = 0.0125 using trigonometric B-spline (Middle) and B-spline (Right).

The comparison of the L_2 and L_{∞} errors is done at different time levels with $\Delta t = 0.0001$ and h = 0.0125with the errors obtained by the earlier schemes: CMMCB²⁶, RBFCM²¹ and QBCM²⁵, and are presented in Table 4 and it can be seen that the obtained results are better than the results given in^{21,25,26}. The physical behaviour of the MCB-DQM and MTCB-DQM solutions at various time levels $t \leq 5$ are shown in Figure 4.

Table 4. Comparison of L_2 and L_{∞} errors at different time levels ≤ 5 with the errors due to well-known earlier schemes

	MTCB - DQN	1	MCB - DQM		QBCM ²⁵
t	h = 0.0125,	$\Delta t = 0.001$	h = 0.0125,	$\Delta t = 0.001$	$h = 0.005, \Delta t = 0.01$
	L_2	L_{∞}	L_2	L_{∞}	L_{∞}
1	4.00E - 05	5.40E - 05	4.20E - 05	5.58E - 05	1.9175E - 04
2	2.51E - 06	5.32E - 06	6.49E - 06	9.71E - 06	1.1387E - 04
3	5.72E - 06	8.48E - 06	3.96E - 06	7.00E - 06	1.7053E - 04
4	1.40E - 05	2.00E - 05	1.19E - 05	1.76E - 05	2.0271E - 04
5	5.26E - 06	7.49E - 06	4.23E - 06	3.01E - 06	9.8405E - 05
	CMMCB ²⁶		RB	FCM ²¹	
t	h = 0.01,	$\Delta t = 0.001$	h = 0.01,	$\Delta t = 0.001$	
	L_2	L_{∞}	L_2	L_{∞}	
1	4.55E - 05	1.43E - 05	1.43E - 04	1.84E - 05	
-				A1016 00	
2	1.43E - 05	8.08E - 05	8.08E - 05	1.07E - 05	
2	1.43E - 05 6.42E - 06	8.08E - 05 1.29E - 05	8.08E - 05 1.29E - 04	1.07E - 05 1.81E - 05	
2 3 4	1.43E - 056.42E - 068.92E - 06	8.08E - 05 1.29E - 05 1.18E - 05	8.08E - 05 1.29E - 04 1.18E - 04	1.07E - 05 1.81E - 05 1.64E - 05	

Example 3: The telegraph equation over the region [0, 2] is considered with conditions.

$$u(0,t) = \tan\left(\frac{t}{2}\right), \qquad u(2,t) = \tan\left(\frac{2+t}{2}\right), \qquad t \ge 0$$

and the function

$$f(x,t) = \alpha \left(1 + \tan^2\left(\frac{x+t}{2}\right)\right) + \beta^2 \tan^2\left(\frac{x+t}{2}\right)$$

The analytical solution¹⁵

$$u(x,t) = \left(\frac{x+t}{2}\right)$$

Numerical solutions computed for are $\alpha = 10, \beta = 5$ taking $\Delta t = 0.001, 0.0001$ and h = 0.025 at diverse time levels $t \le 1$. The comparison of L_{∞} error is done at different time levels with earlier schemes^{23,25,26} and are reported in Table 5. It is evident from Table 5 and 6 that obtained results are matched with the analytical solution and also better than the previously obtained results. The behaviour of the MCB-DQM and MTCB-DQM numerical solutions are depicted physically in Figure 4 at different time level $t \leq 1.0$, taking $\Delta t = 0.0001$ and h = 0.025 h = 0.025



Figure 4. At different time levels $t \le 1$ the performances of exact (Left) and approximate solution of Example 3 with $\Delta t = 0.0001$ and h = 0.025 using trigonometric B-spline (Middle) and B-spline (Right).

5. Conclusion

Solution of the hyperbolic telegraph equation is obtained numerically in this paper using two different schemes named as, MCB-DQM and MTCB-DQM. These schemes are based on the DQM combined with modified cubic B-spline and modified cubic trigonometric B-spline as basic functions. On implementing the schemes and substituting the derivatives, set of ODEs are attained, which is solved using SSP RK43. The efficiency and precision of the proposed method is revealed by three test problems. The numerical result, L_2 and L_{∞} errors are compared with numerical solutions from literature and are found to be in decent agreement with formerly obtained results. The advantage of the developed methods is there ease to implement and reduced data complexity as compared to the present schemes.

6. References

- El-Azab MS, El-Ghamel A. Numerical Algorithm for the Solution of Telegraph Equations, Applied Mathematics Computing. 2007; 190(1):757–64.
- Mohanty RK. New Unconditionally Stable Difference Schemes for the Solution of Multidimensional Telegraphic Equations, International Journal of Computing Mathematics. 2009; 86(12):2061–71.
- 3. Jordan PM, Puri A. Digital Signal Propagation in Dispersive Media, Journal of Applied Physics. 1999; 85(3):1273–82.
- Weston VH, He S. Wave Splitting of the Telegraph Equation in R³ and its Application to Inverse Scattering, Inverse Problems. 1993; 9(6):789–812.
- 5. Banasiak J, Mika JR. Singularly Perturved Telegraph Equations with Applications in the Random Walk Theory,

Schemes	$(h, \Delta t)$	t = 0.2	t = 0.4	t = 0.6	t = 0.8	t = 1.0
MCB - DQM	(0.025, 0.0001)	4.85E-05	4.91E-05	2.26E-05	4.33E-04	1.12E-02
MTCB – DQM	(0.025, 0.0001)	2.42E-04	3.80E-04	6.98E-04	1.73E-03	1.64E-02
CMMCB ²⁶	(0.020, 0.0001)	3.47E-05	5.34E-05	9.47E-05	1.87E-04	5.87E-04
MCB – DQM	(0.025, 0.001)	9.62E-04	1.47E-03	2.56E-03	5.61E-03	1.84E-02
MTCB – DQM	(0.025, 0.001)	2.42E-03	3.80E-03	6.97E-03	1.72E-02	9.92E-02
CMMCB ²⁶	(0.020, 0.001)	2.63E-04	6.99E-04	1.48E-03	3.40E-03	1.11E-02
<i>CBQ</i> ²³	(0.005, 0.001)	1.89E-04	3.99E-04	7.97E-04	1.87E-03	8.01E-03
QBCM ²⁵	(0.001, 0.001)	2.77E-04	7.07E-04	1.38E-03	3.09E-03	1.34E-02

Table 5. Comparison of L_{∞} errors in Example 3 at different time levels $t \leq 1$ with the errors in the earlier schemes

Table 6. Comparison of L_2 errors in Example 3 at different time levels $t \le 1$ with the errors in the earlier scheme²⁶

Schemes	$(h, \Delta t)$	t = 0.2	t = 0.4	t = 0.6	t = 0.8	t = 1.0
MCB – DQM	(0.025, 0.0001)	1.31E-05	1.48E-05	1.01E-05	8.56E-05	1.93E-03
MTCB – DQM	(0.025, 0.0001)	4.85E-05	8.18E-05	1.40E-04	3.30E-04	3.33E-03
CMMCB ²⁶	(0.020, 0.0001)	5.03E-05	9.51E-05	2.20E-04	7.82E-04	7.92E-03
MCB – DQM	(0.025, 0.001)	2.41E-04	3.82E-04	6.45E-04	1.31E-03	3.97E-03
MTCB – DQM	(0.025, 0.001)	5.82E-04	1.02E-03	1.82E-03	3.95E-03	1.76E-02
CMMCB ²⁶	(0.020, 0.001)	1.87E-04	4.88E-04	9.48E-04	1.87E-03	5.09E-03

Journal of Applied Mathematics Stochastic Analysis. 1998; 11(1):9-28.

- 6. Industrial Microwave. Date Accessed: 2012. http://www. industrialmicrowave.com/.
- 7. Roussy G, Pearcy JA. Foundations and Industrial Applications of Microwaves and Radio Frequency fields. Wiley: New York, 1995.
- Kim WS. Doubly-Periodic Boundary Value Problem for Nonlinear Dissipative Hyperbolic Equations, Journal of Mathematical Analysis and Application. 1990; 145(1):1-16.
- 9. Mohanty RK. An Unconditionally Stable Difference Scheme for the one Space–Dimensionallinear Hyperbolic Equation, Applied Mathematics Letter. 2004; 17(1):101–5.
- Mohanty RK. An Unconditionally Stable Finite Difference Formula for a Linear Second Order one Space Dimensional Hyperbolic Equation with Variable Coefficients, Applied Mathematics Computing. 2005; 165(1):229-36.
- 11. Liu HW, Liu LB. An Unconditionally Stable Spline Difference Scheme of $O(k^2 + h^2)$ for Solving the Second

Order 1D Linear Hyperbolic Equation, Mathematical Computation Modelling. 2009; 49(9-10):1985–93.

- Dehghan M. On the Solution of an Initial–Boundary Value Problem that Combines Neumannand Integral Condition for the Wave Equation, Numerical Methods Partial Differential Equations. 2005; 21(1):24–40.
- 13. Mohanty RK, Jain MK, George K. On the use of High Order Difference Methods for the system of One Space Second Order Non-Linear Hyperbolic Equations with Variable Coefficients, Journal of Computational Applied Mathematics. 1996; 72(2):421-31.
- Twizell EH. An Explicit Difference Method for the Wave Equation with Extended Stability Range, BIT Numerical Mathematics. 1979; 19(3):378–83.
- Mohebbi A, Dehghan M. High Order Compact Solution of the One-Space–Dimensional Linear Hyperbolic Equation, Numerical Methods Partial Differential Equations. 2008; 24(5):1222-35.
- Dehghan M, Lakestani M. The use of Chebyshev Cardinal Functions for Solution of the Second-Order One– Dimensional Telegraph Equation, Numerical Methods Partial Differential Equations. 2009; 25(4):931-38.
- Lakestani M, Saray BN. Numerical Solution of Telegraph Equation using Interpolating Scaling Functions, Computational Mathematics Application. 2010; 60(7):1964–72.
- Saadatmandi A, Dehghan M. Numerical Solution of Hyperbolic Telegraph Equation using the Chebyshev Tau Method, Numerical Methods Partial Differential Equations. 2010; 26(1):239–52.
- Gao F, Chi CM, Unconditionally Stable Difference Schemes for a One Space– Dimensional Linear Hyperbolic Equation, Applied Mathematics Computation. 2007; 187(2):1272–76.
- Dehghan M, Ghesmati A. Solution of the Second-Order One–Dimensional Hyperbolic Telegraph Equation by using the Dual Reciprocity Boundary Integral Equation (DRBIE) Method, Engineering Analytical Boundary Element. 2010; 34(1):51–59.
- 21. Dehghan M, Shokri A. A Numerical Method for Solving the Hyperbolic Telegraph Equation, Numerical Methods Partial Differential Equations. 2008; 24(4):1080–93.
- 22. Jiwari R, Pandit S, Mittal RC. A Differential Quadrature Algorithm for the Numerical Solution of the Second-Order One Dimensional Hyperbolic Telegraph Equation,

International Journal of Nonlinear Science. 2012; 13(3):259–66.

- 23. Dosti M, Nazemi A. Solving One–Dimensional Hyperbolic Telegraph Equation using Cubic Bsplinequasi– Interpolation, World Academy of Science, Engineering and Technology. 2011; 5(4):935–40.
- Rashidinia J, Jamalzadeh S, Esfahani F. Numerical Solution of One Dimensional Telegraph Equation using Cubic B-Spline Collocation Method, Journal of Interpolation Approximation in Scientific Computing. 2014(2014); 1–8.
- 25. Dosti M, Nazemi A. Quartic B-Spline Collocation Method for Solving One–Dimensional Hyperbolic Telegraph Equation, Journal of Information Computing Science. 2012; 7(2):83–90.
- Mittal RC, Bhatia R. Numerical Solutions of Second Order One Dimensonal Hyperbolic Telegraph Equation by Cubic B-Spline Collocation Method, Applied Mathematics and Computation. 2013; 220(1):496–506.
- 27. Arora G, Singh BK. Numerical Solution of Burgers Equation with Modified Cubic B-Spline differential Quadrature Method, Applied Mathematics Computation. 2013; 224(1):166–77.
- Spiteri JR, Ruuth SJ. A New Class of Optimal High-Order Strong Stability-Preserving Time stepping Schemes; SIAM Journal Numerical Analysis. 2002; 40(2):469–91.
- Bellman R, Kashef BG, Casti J. Differential Quadrature: A Technique for the Rapid Solution of Nonlinear Differential Equations, Journal of Computational Physics. 1972; 10(1):40–52.
- Korkmaz A. Shock Wave Simulations using Sinc Differential Quadrature Method. International, Journal of Computational Aided Engineering Software. 2011; 28(6):654–74.
- 31. Quan JR, Chang CT. New Insights in Solving Distributed System Equations by the Quadrature Methods–I; Computers and Chemistry Engineering. 1989; 13(9):779–88.
- 32. Shu C, Richards BE. Application of Generalized Differential Quadrature to Solve Two Dimensional in compressiblenavier-Stokes Equations, International Journal of Numerical Methods in Fluids. 1992; 15(7);791–8.
- 33. Shu C. Differential Quadrature and its Application in Engineering, Springer-Verlag London Limited, 2000.